



# Long term variability of the solar dynamo

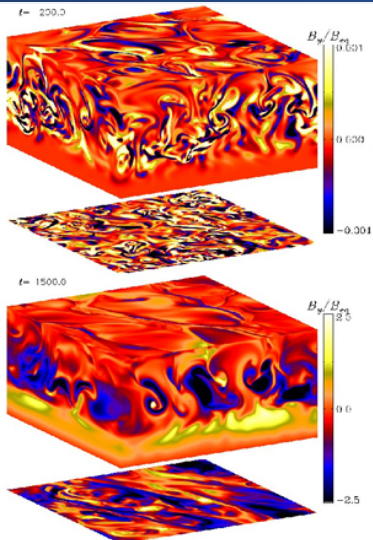
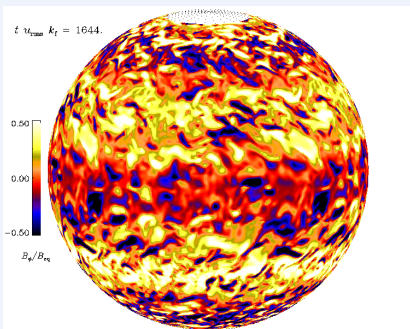
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Jörn Warnecke<sup>2</sup>, Axel Brandenburg<sup>3</sup>

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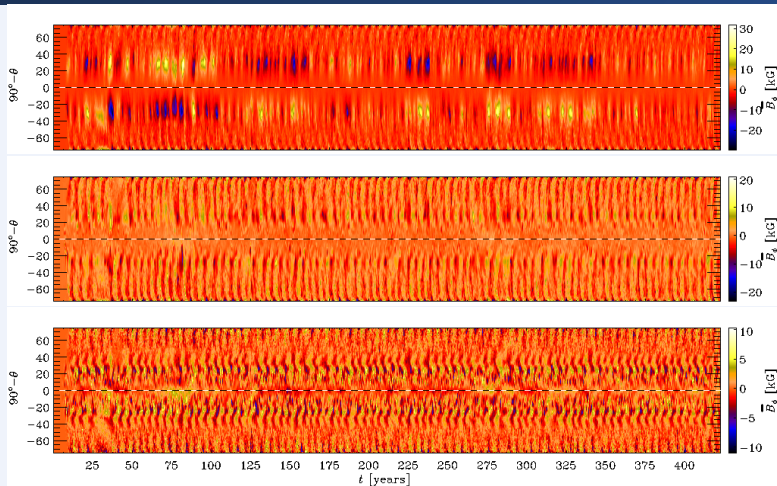
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Global spherical convection dynamo e.g. (Käpylä et al. 2013, Käpylä et al. 2015)



**Figure:**  $B_\phi$  averaged azimuthally as function of latitude over time - layers near the base, middle and surface of the convection zone. Time derived by  $5\Omega_\odot / R_\odot$ , for a solar size star rotating 5x solar rate

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$$\mathbf{F}^{\text{rad}} = -K \nabla T \quad \text{and} \quad \mathbf{F}^{\text{SGS}} = -\chi_{\text{SGS}} \rho T \nabla s \quad (5)$$

are heat fluxes, radiative and SGS (sub grid scale - numerical stability)



<b><math>A</math></b>	magnetic vector potential
<b><math>U</math></b>	velocity
<b><math>B = \nabla \times A</math></b>	magnetic field
<b><math>J = \mu_0^{-1} \nabla \times B</math></b>	current density
$\mu_0$	vacuum permeability
<b><math>D/Dt = \partial/\partial t + \mathbf{u} \cdot \nabla</math></b>	material derivative
<b><math>S</math></b>	rate of strain tensor
$\rho$	density
$\nu$	kinematic viscosity
$\eta$	magnetic diffusivity
$K$	radiative heat conductivity
$\chi_{SGS}$	turbulent heat conductivity (unresolved convective transport of heat)
$s$	specific entropy
$T$	temperature
$p$	pressure

Ideal gas law:  $p = (c_P - c_V)\rho T$ , where adiabatic index  $\gamma = c_P/c_V = 5/3$ .

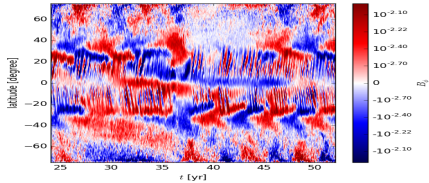


Figure:  $\langle B_\phi \rangle_\phi$  near surface of the convection zone during grand minima south then north.

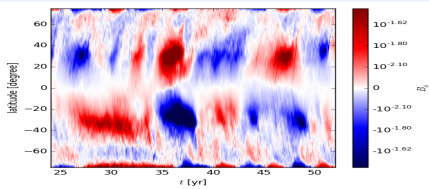


Figure:  $\langle B_\phi \rangle_\phi$  near base of the convection zone during grand minima south then north.

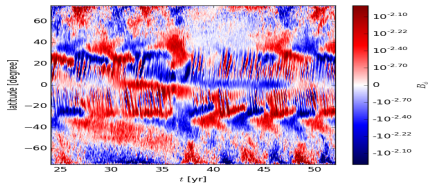


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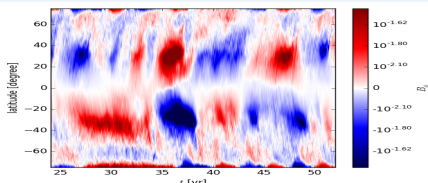


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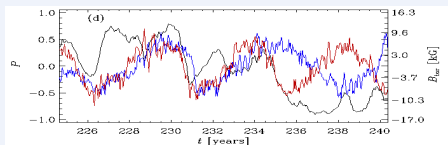


Figure: Parity (black) and  $\langle B_\phi \rangle_\phi$  near surface (N:blue, S:red) at  $\pm 25^\circ$  latitude, during high state of base toroidal mode

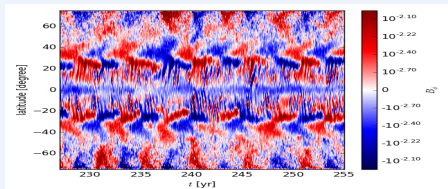


Figure:  $\langle B_\phi \rangle_\phi$  near the surface of the convection zone during switch from N-S symmetry to asymmetry.

(Schrinner et al. 2007) expressed the induction equation in terms of the mean field (e.g. azimuthal average)

such that  $\mathbf{B} = \overline{\mathbf{B}} + \mathbf{b}$  and  $\mathbf{U} = \overline{\mathbf{U}} + \mathbf{u}$

Taking the curl of (1)

$$\frac{\partial}{\partial t}(\overline{\mathbf{B}} + \mathbf{b}) = \nabla \times (\overline{\mathbf{U}} + \mathbf{u}) \times (\overline{\mathbf{B}} + \mathbf{b}) + \eta \nabla^2 (\overline{\mathbf{B}} + \mathbf{b}), \quad (6)$$

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$$\frac{\partial \overline{\mathbf{B}}}{\partial t} = \nabla \times (\overline{\mathbf{U}} \times \overline{\mathbf{B}}) + \nabla \times \mathcal{E} + \eta \nabla^2 \overline{\mathbf{B}}, \quad (7)$$

where the electromotive force (EMF)  $\mathcal{E} = \overline{\mathbf{u} \times \mathbf{b}}$

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where  $\mathbf{G} = \mathbf{u} \times \mathbf{b} - \overline{\mathbf{u} \times \mathbf{b}}$

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Further assume only weakly in space and time

In general, if no higher than first order spatial derivatives and no time derivatives of  $\bar{\mathbf{B}}$  are taken into account then in general

$$\mathcal{E} = \mathbf{a}\bar{\mathbf{B}} + \mathbf{b}\nabla\bar{\mathbf{B}}$$

with second rank  $\mathbf{a}$  and third rank  $\mathbf{b}$  tensors  
i.e. 36 independent coefficients



in curvilinear coordinates can be expressed

$$\mathcal{E} = \alpha \bar{\mathbf{B}} + \gamma \times \bar{\mathbf{B}} - \beta \cdot (\nabla \times \bar{\mathbf{B}}) - \delta \times (\nabla \times \bar{\mathbf{B}}) - \kappa \cdot (\nabla \bar{\mathbf{B}})^{\text{sym}} \quad (9)$$

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Coefficients, vectors  $\gamma$  and  $\delta$ , second rank tensors  $\alpha$  and  $\beta$ , and third rank tensor  $\kappa$  represent a decomposition of the velocity field and can be related to physical processes, helping us to understand the structure of the dynamo.

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How might we determine these coefficients?

Assume Eq.(8) for  $\mathbf{b}$  pertains to a steady test field  $\overline{\mathbf{B}}_T$

$$\frac{\partial \mathbf{b}}{\partial t} - \nabla \times (\overline{\mathbf{U}} \times \mathbf{b}) - \nabla \times \mathbf{G} - \eta \nabla^2 \mathbf{b} = \nabla \times (\mathbf{u} \times \overline{\mathbf{B}}_T) \quad (10)$$

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For any given test field  $\overline{\mathbf{B}}_T^{(i)}$  the EMF, depending only on  $\overline{\mathbf{U}}$  and  $\mathbf{u}$ , can be expressed

$$\mathcal{E}^{(i)} = \tilde{a}_{jk} \overline{\mathbf{B}}_{T_k}^{(i)} + \tilde{b}_{jkr} \frac{\partial \overline{\mathbf{B}}_{T_k}^{(i)}}{\partial r} + \tilde{b}_{jk\theta} \frac{1}{r} \frac{\partial \overline{\mathbf{B}}_{T_k}^{(i)}}{\partial \theta} \quad (11)$$

Applying these to 9 linearly independent axisymmetric test fields we can solve simultaneously for the 27 independent coefficients

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How plausible is the mean field model?

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It depends on the extent to which the EMF satisfies the assumptions



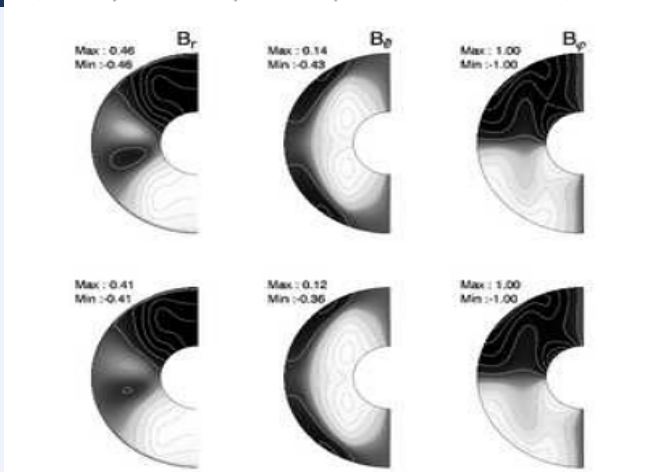
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How does the mean field dynamo compare with the full solution?

*M. Schrinner, K.-H. Rädler, D. Schmitt, M. Rheinhardt and U. R. Christensen*



**Figure:** (Schrinner et al. 2007) magnetoconvection: azimuthally averaged magnetic field components resulting from DNS (upper), mean-field calculations derived from test field (lower).  $[(\rho\mu_0\eta\Omega)^{1/2}]$

*M. Schurrner, K.-H. Rädler, D. Schmitt, M. Rheinhardt and U. R. Christensen*

Max: 2.66  
Min: -4.08



Max: 2.16  
Min: -2.16



Max: 2.34  
Min: -0.80



Max: 2.55  
Min: -4.24



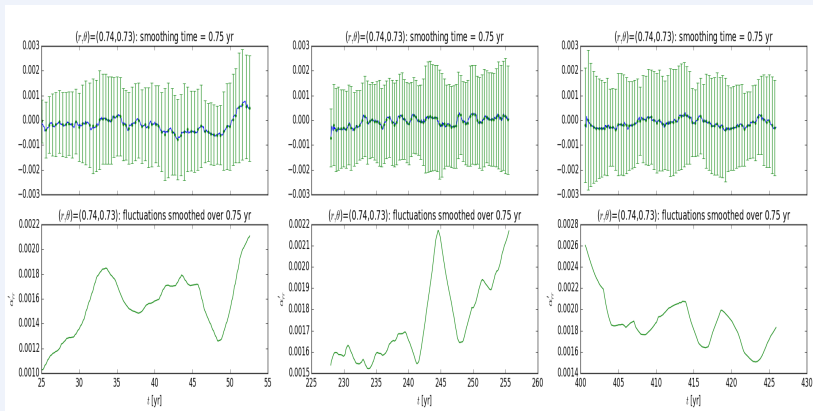
Max: 2.04  
Min: -2.04



Max: 1.87  
Min: -0.59



**Figure:** (Schurrner et al. 2007) electromotive forces in the magnetosphere (top)  $\mathcal{E}_r^{\text{MHD}}$ ,  $\mathcal{E}_\theta^{\text{MHD}}$ ,  $\mathcal{E}_\phi^{\text{MHD}}$ , and (bottom)  $\mathcal{E}_r^{\text{MF}}$ ,  $\mathcal{E}_\theta^{\text{MF}}$ ,  $\mathcal{E}_\phi^{\text{MF}}$ .  
 $[(\eta/D)(\rho\mu_0\eta\Omega)^{1/2}]$



**Figure:** Upper:  $\alpha_{rr}(0.74R_{\odot}, 0.73\text{rad})$ , box-car averaged over 9 month intervals (blue), standard deviation over interval (error bars). Lower: Perturbations  $\alpha'_{rr_{rms}}$  box-car averaged over 9 months

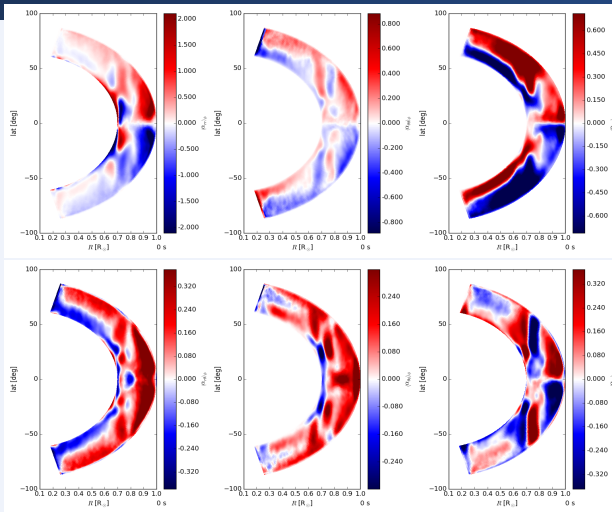
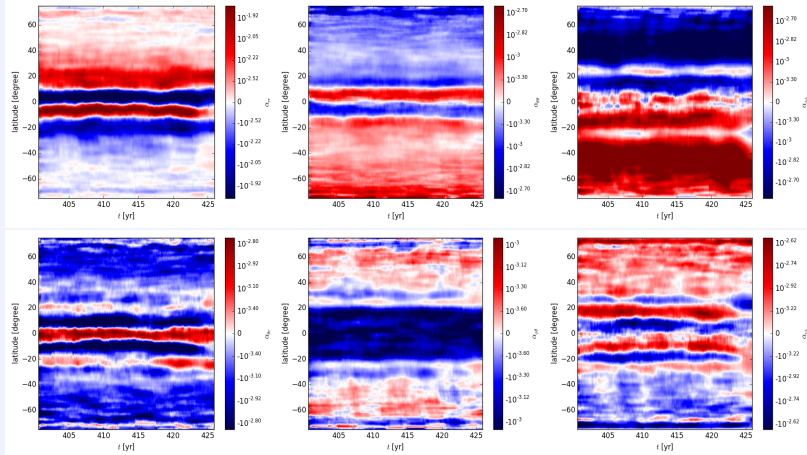
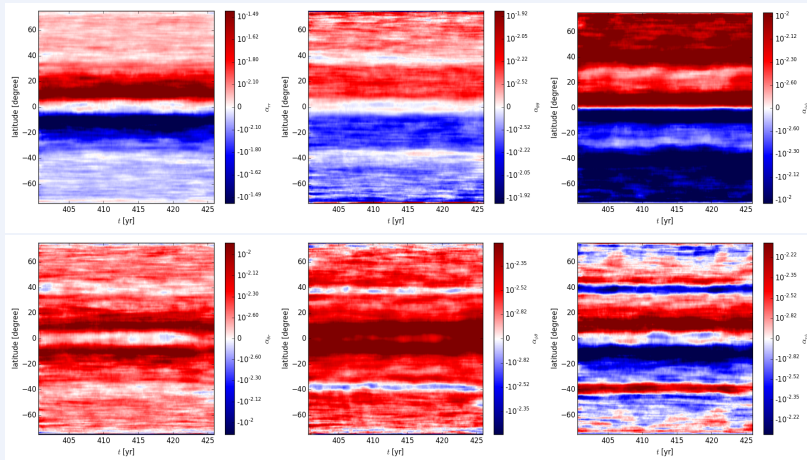


Figure: Time and azimuthally averaged  $\alpha$ -tensor for the period corresponding to steady cyclic epoch



**Figure:** Time evolution of azimuthally averaged  $\alpha$ -tensor spanning steady cyclic epoch near the base of the convection zone



**Figure:** Time evolution of azimuthally averaged  $\alpha$ -tensor spanning steady cyclic epoch near the surface of the convection zone

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- ▶ mean field model may work with simple random  
fluctuations?

- Käpylä M J, Käpylä P J, Olsper N, Brandenburg A, Warnecke J, Karak B B & Pelt J 2015 *ArXiv e-prints* .
- Käpylä P J, Mantere M J, Cole E, Warnecke J & Brandenburg A 2013 *ApJ* **778**, 41.
- Schrinner M, Rädler K H, Schmitt D, Rheinhardt M & Christensen U R 2007 *Geophysical and Astrophysical Fluid Dynamics* **101**, 81–116.